### The tale of angular momentum transport in primordial galaxies and the formation of massive BH seeds

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### The puzzle of high-z Quasars

Bright Quasars (L > 10<sup>47</sup> erg/s) < 700 million years after Big Bang (z > 6) M<sub>BH</sub> > 10<sup>9</sup> Mo (see X. Fan's talk f= Mortlock et al. 2011; Banados et al. 2016; 2019)  $M(t) = M e^{\varepsilon} e^{\int_{Edd} 0.45 Gyr}$ PROBLEM: is there enough time to grow these early SMBHs?



High-z QSO rare (~10<sup>-8</sup> h<sup>3</sup> Mpc<sup>-3</sup>), 4 orders of magnitude less abundant than their z=0 counterparts. Abundance and clustering suggest their hosts rare massive halos ( $M_{halo} > 10^{12}$  Mo – see Volonteri & Rees 2006; Sijacki et al. 2010)

High-z QSOs hardly from Eddington-limited accretion of Pop III BH seeds as seed mass should be >~ 10<sup>4</sup> Mo Alternative scenarios:

(1) Direct gas collapse into black hole seeds (10<sup>4</sup>-10<sup>5</sup> Mo) (eg Begelman et al. 2006; Lodato & Natarajan 2006; Begelman 2010; Mayer et al. 2010; Latif et al. 2014;2015; Regan & Haenhelt. 2009; 2015; Mayer & Bonoli 2019; Woods et al. 2019; Wyse et al. 2019)



Formation of massive BH from collapse of massive gas cloud, with a prior stage as a supermassive star (SMS) and/or quasi-star (Begelman 2010), or even directly into a BH via GR radial instability ("dark collapse" see Mayer & Bonoli 2019; Hammerle et al. 2019)

(2) Super-Eddington accretion onto Pop III seeds (eg Madau et al. 2014; Lupi et al (2016); Inayoshi et al. 2016; Regan et al. 2019, see review in Mayer 2018, 2018arXiv180706243M)

Regime: radiatively inefficient accretion in large optical depth medium ( $\tau >~ 10^3$ ) with highly viscous-driven accretion flow tdiff >> tadvection Prototypical model is SLIM disk by Sadowski et al. (2011;2013). Low radiative efficiency (~ 0.1%, as opposed to >~10% in standard α-disk accretion), no energy loss via winds

More recent 3D MHD-RT accretion disk simulations by Jiang et al. (2015;2018) using VET method for RT show enhanced radiative losses in turbulent non-axisymmetric flow with **radiative efficiency** of 1-5%, sustaining M<sub>dot</sub> ~ 25-150 M<sub>dot</sub>, edd. (+winds) In both direct gas collapse and Super-Eddington accretion a **challenge is to supply the gas at sufficiently high rates to sub-nuclear scales (< 10<sup>-2</sup> - 1 pc — scales of supermassive protostellar cloud in direct collapse scenarios or the accretion disk in super-Eddington scenarios)** 

For example, for SMS formation ( $M^* > 10^4$  Mo)  $M_{dot} > 0.1-1$  Mo/yr is required (Hosokawa et al. 2014; Woods et al. 2019). For Super-Eddington accretion onto light Pop III BH seeds (<1000 Mo) (Madau & Rees 2000) smaller M<sub>dot</sub> required at the beginning then larger as seed grows (Regan et al. 2019) For dark collapse into a massive BH even larger  $M_{dot} >~ 100$  Mo/yr is required (Hammerle, Meynet, Mayer et al. 2019).

Need to feed from gas reservoir in galactic-scale disk Which angular momentum transport mechanism(s) can feed the sub-nuclear galactic region efficiently enough?

### A GENERAL MECHANISM FOR ANGULAR MOMENTUM TRANSPORT: GRAVITATIONAL TORQUES

Toomre parameter  $Q = \kappa c_s / \pi G \Sigma$ 

(balance centrifugal force, pressure and gravity)



#### THREE REGIMES

(a) Q <~ 1 locally unstable to collapse, fragmentation (eg star formation) on a dynamical timescale (t<sub>dyn</sub>)

(b) 1 < Q < 2 locally stable, globally unstable to non-axisymmetric disturbances (spiral modes, bar-like modes)</li>
Angular momentum transport (on ~ t<sub>orb</sub> timescale)
via spiral density waves (Lynden Bell & Pringle1979; Lin & Pringle 1987; Laughlin & Adams 2000)
Mass/angular momentum transport can be parametrized with (local) 'alpha disk"
(effective "alpha" large <~ 0.1).</li>
--> Nonlinear regime, hydro simulations needed

(c) Q > 2 locally and globally stable - dynamically uninteresting (remains close to axisymmetric)

### 1 < Q < 2 is the "useful" regime for sustained angular momentum transport

The cooling timescale controls in which Q regime the disk will settle. IF  $t_{cool} < t_{orb} \rightarrow Q < 1 \rightarrow fragmentation$  (Gammie 2001;Rice et al. 2004; Deng et al. 2017)

How can one enforce  $t_{cool} > t_{orb}$ ? Conventional route: inefficient radiative cooling by ansatz — metal-free gas plus dissociation of H<sub>2</sub> byLyman-Werner ionising radiation field from nearby star forming galaxies—-> atomic cooling halos at z ~ 15-20 (eg Wyse et al. 2019)



Example of protogalactic disk simulation in metal-free stomic cooling halo Latif et al. (2013)

Jeans unstable clump (M >~ 10<sup>4</sup> Mo) M<sub>dot</sub> <~ 1 Mo/yr Mclump >~ 10<sup>3-4</sup> Mo
(still too small BH seed?) Alternative route direct collapse route: nuclear inflows in mergers of massive metal-enriched galaxies at z ~ 8-10 (Mayer et al. 2010;2015; Mayer & Bonoli 2019) (M<sub>vir</sub> ~ 10<sup>12</sup> Mo, 4-5σ peaks)

> The inner 200 pc region a few Myr before final merger: the remnants of the two galaxy cores are shown



Parsec- scale dense nuclear disk forms which keeps accreting matter infalling supersonically (Mdot > 1000 Mo/yr). Disk core is warm and stable to fragmentation due to dynamical heating by shocks in accretion flow and disk (Mayer et al. 2015)

>~10<sup>9</sup> Mo inside ~2 pc in only ~ 10<sup>4</sup> yr after galaxy merger Any direct collapse pathway possible, including "dark collapse" (Hammerle et al. 2019) A novel regime: Magnetised Gravitoturbulent Nuclear Disk Increased stability against fragmentation? New pathway for enhanced angular momentum transport?



Test case: self-gravitating protoplanetary disk

Assume some prior amplification of B field by turbulent dynamo during gravitational collapse or merger (eg Schober et al. 2013; Grete et al. 2019) —> Initial seed B<sub>z</sub> field in the range 0.1-1 Gauss

Regime initially studied in 3D magnetised self-gravitating protoplanetary disk simulations (Deng, Mayer et al. 2019)

## Gravitational instability with magnetic dynamo: stabilising effect (Q rises) PLUS enhanced turbulent mass transport



#### Hydro

MHD

Disk models with a range of cooling rates, from a few to 10 orbital times

 $\alpha \sim \langle H_{r\phi} + M_{r\phi} + G_{r\phi} \rangle / \langle P \rangle$ 

Stresses (hydro, magnetic and gravitational)

Turbulent viscosity a ~0.2-0.3 In non-magnetised self-gravitating disks a <~ 0.1

### Shown is vertical slice through disk

### Field amplification driven by dynamo mechanism

Dynamo generated by vertical circulation around spiral density wave intensity maxima



Magnetic energy grows >~10 times larger than in identical. non-self gravitating disk with MRI (plasma  $\beta$  ~ 7-10 rather than  $\beta$  >~ 100)

# First local shearing box simulations by **Riols and Latter 2017;2018**) with PLUTO code uncovered gravitational instability (GI) driven spiral dynamo loop:

(I) radial compression of B field by spiral density waves + (II) lifting and folding of B field by vertical rolls generates new radial field + (III) shearing of new radial field back into toroidal by differential rotation (toroidal -> poloidal -> toroidal loop with net mean field growth)



Preliminary results: field amplification in self-gravitating nuclear disk in high-z merger remnant in Mayer et al. (2015)



()Caution: current simulation *adiabatic* without cooling gravitational instability weakens over time —-> weaker dynamo trigger
 () Will need to add cooling to assess increased stability against fragmentation